

A note on Williams decomposition of the BES(3) process (19 Nov. 2012)

Abstract: We present a concise proof of Williams decomposition of the BES(3) process, starting from  $r > 0$ , at its ultimate minimum.

The proof is done with the help of the progressive enlargement formula with respect to  $\mathbb{F}_t$ , the time of the ultimate minimum.

This discussion is strongly motivated by, and linked with, the reading guide of Mikogahara-Peter about general fast passage times.

### 1. Introduction and Some Statements from [NP].

(1.1) In their survey [NP] about fast passage times, the authors illustrate some of their formulas with the following example:

Let  $(R_t, t \geq 0)$  be a BES(3) process on  $\mathbb{R}_+$ , starting from  $r > 0$ . Denote by  $(\tilde{R}_t)_{t \geq 0}$  its natural filtration,

and let  $I_t = \inf_{s \leq t} R_s$ .

The following results are found in [NP], around Corollary 4.10

Appendix: The authors are grateful to Monique Jeanblanc for providing them with a preprint of [NP].

c)  $T_\infty$  is distributed as  $nU$ , with  $U$  uniform on  $(0,1)$ .

b) The (Azéma-) supermartingale associated with  $g$ , the time at which  $R$  reaches  $T_\infty$  is:

$$Z_t \equiv P(g > t | \mathcal{R}_t) = I_t / R_t.$$

c) The Laplace transform of the law of  $g$  is:

$$E[\exp(-\lambda g)] = \frac{1}{1 + \sqrt{2\lambda} n} \{ 1 - \exp(-\sqrt{2\lambda} n) \}$$

d) The density of  $g$  is:

$$h(t) = \frac{1}{1 + \sqrt{2\lambda} n} (1 - e^{-\sqrt{2\lambda} n})$$

(1.2)

Recall that, if  $(B_t, t \geq 0)$  is a Brownian motion starting from 0, and  $a \in \mathbb{R}$ , then the law of the first hitting time of  $a$  by  $B$ , denoted as  $T_a^{(B)}$  is:

$$(1) \quad P(T_a^{(B)} \in dt) = \frac{\sqrt{2\pi t^3}}{dt} |a| \exp(-a^2/2t).$$

This well known fact allows to rewrite c) and d) above as follows:

$$(2) \quad g \stackrel{\text{law}}{=} T_{nU}^{(B)}$$

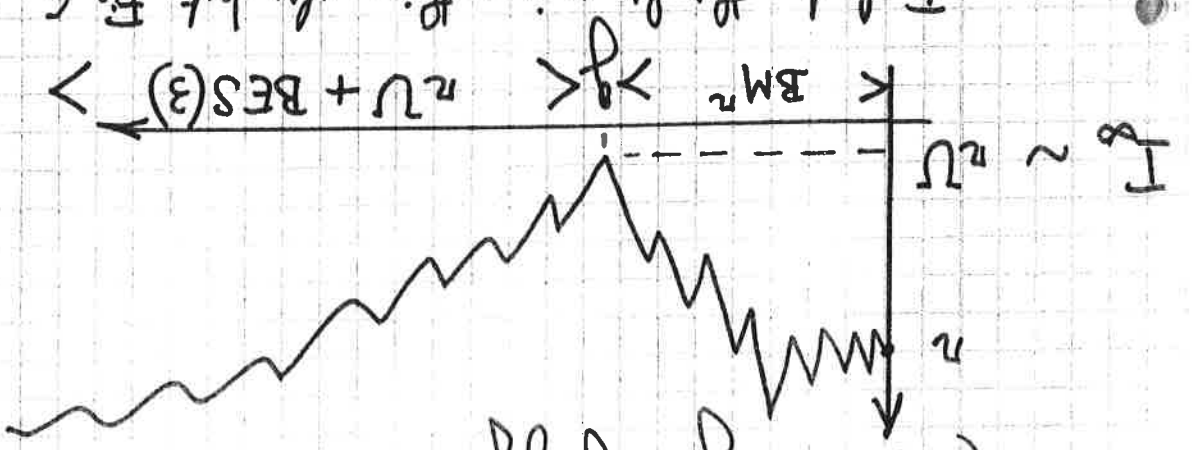
where  $U$ , uniform on  $[0,1]$ , is independent from  $B$ .

We let the reader verify (2) from c) and d).

In fact, (2) may be understood via the classical decomposition of  $R$  before and after  $g$ , due to Williams [W]

2. Williams decomposition of  $R_t$ , before and after  $g$ , via programic enlargement.

(2.1) The following figure divides with the decomposition.



In fact, the figure is nothing else but Fig 5 in Revuz-Yor ([RY], Proposition 3.10 and Theorem 3.11 in Chap 6, Sect 3) when the BES(3) path is considered starting from level  $c (= r, \text{ here})$ .

(2.2) We now state precisely Williams' decomposition theorem before and after  $g$ .

Theorem (Williams [W]) Consider the following 3 independent random objects:

- (i)  $(B_t, t \geq 0)$ , a BM starting from  $r > 0$ ;
- (ii)  $U$  a uniform r.v. on  $[0, 1]$
- (iii)  $(\tilde{R}_t, t \geq 0)$  a BES(3) process starting from 0.

then, the process:

$$R'_t, \text{ for } t \leq g \equiv \inf\{u: B'_u = 0\}$$

$$R_t =$$

(3)  $rU + \tilde{R}(t-g)$ , for  $t > g$

is a BES(3) process starting from  $r$ .

We note that the pre-g-Brownian motion found in (3) explains the result (2). Indeed, if  $B_t^{(0)} = n - B_t^{(0)}$ , then: (2')  $g = \inf\{u : B_u^{(0)} = n(1-u)\}$  from which (2) follows.

(2.3) We now proceed to the proof of the theorem, via the enlargement formula which describes the additive decomposition of the BES(3) process  $(R_t)$  in the filtration  $(\mathcal{R}_t^g)$  containing  $(R_t)$ , and making of a stopping time —

We have:

(4)  $R_t = n + B_t + \int_0^t \frac{d\alpha}{R_\alpha}$ , with  $(B_t)$  a Brownian motion with respect to  $(\mathcal{R}_t)$ . (see, e.g., [J])

Then, the enlargement formula in  $(\mathcal{R}_t^g)$  yields:

(5)  $n + B_t = B_t' + \int_0^t \frac{d\alpha}{R_\alpha} + \int_0^g \frac{z_u}{d < B, z >_u} + \int_0^g \frac{(1-z_u)}{d < B, 1-z >_u}$  with  $(B_t')$  a Brownian motion with respect to  $(\mathcal{R}_t^g)$ .

Moreover, we deduce from b) that:

(6)  $d < B, z >_u = -\frac{z_u}{dR_u}$ ,  $u \leq g$

(7)  $\frac{1-z_u}{d < B, 1-z >_u} = \frac{I_{\infty, dR_u}(R_u - I_{\infty})}{R_u}$ ,  $u > g$

These two identities simply say:

— using (4) and (6), the form of the pre-g-process;

— using (4) and (7), the form of the pr-g-process.

For the proof of (3) to be complete, it remains to prove that  $B'$  is independent of  $I_\infty$  (law  $\nu$ ), or, more exactly that, given  $I_\infty = a$ , the pr-g-process in  $(B'_u, u \leq T'_a)$  with obvious notation. This property is asserted in the following Proposition:

Proposition:

Let  $P_n$  denote the law of  $R$ , starting from  $n$ , and let  $P'_n$  denote the law of Brownian motion starting from  $n$ . Then, one has, for  $a < n$ :

$$E_n [F_g | I_\infty = a] \stackrel{(1)}{=} E_n [F_{T_a} | I_\infty < a] \stackrel{(2)}{=} E_n [F_{T'_a} | I_\infty < a] \stackrel{(3)}{=} E_n [F(B'_u, u \leq T'_a)]$$

(8)

Proof of Proposition:

(1) The equality between the RHS of (1) and (3) follows from the classical Doob-absolute continuity relationship:

$$P'_n | F_t = \left( \frac{X_{t-a}}{n} \right) \cdot P_n | F_t$$

(Xt) denotes the coordinate process on path-space on canonical path-space  $C([0, \infty), \mathbb{R})$ . The equality may be extended when replacing  $t$  by a stopping time. In particular

$$P'_n | F_{T'_a} \cap (T_a < \infty) = \left( \frac{a}{n} \right) \cdot P_n | F_{T'_a}$$

we have:

for  $0 < a < n$ , which yields the desired result.

(ii) The equality follows simply from:  $(T < \infty) + (T < \infty) = (T < \infty)$

(iii) It remains to prove (i). We first start with the identity: (for simplicity, we note  $\mathbb{P}$  instead of  $\mathbb{P}_n$ )

$$(9) \quad E[1_{(q \leq t)} \varphi(I_s)] = E\left[\int_t^\infty \varphi(I_s) d(1 - Z_s)\right]$$

for any  $\varphi: [0, \infty) \rightarrow \mathbb{R}_+, \text{Borel}$ . Indeed, the LHS of (9) equals:

$$E[1_{(q \leq t)} \varphi(I_t)] = E[(1 - Z_t) \varphi(I_t)]$$

$$= E\left[\int_t^\infty \varphi(I_s) dI_s (1 - Z_s)\right] + E\left[\int_t^\infty \varphi(I_s) d(1 - Z_s)\right]$$

assuming here that  $\varphi$  is  $C^1$ . We note that the expectation involving  $\varphi'$  is equal to 0, since  $1 - Z_s = 0, dI_s \text{ a.e.}$  Thus, a monotone class argument yields that (9) holds for every  $\varphi \geq 0, \text{Borel}$ .

Next, from the additive decomposition of  $(1 - Z_s)$ , we obtain:

$$E[1_{(q \leq t)} \varphi(I_s)] = E\left[\int_t^\infty \varphi(I_s) \left(-\frac{dI_s}{I_s}\right)\right]$$

$$= E\left[\int_t^\infty \varphi(a) \frac{da}{a}\right]$$

$$= E\left[\int_t^\infty \varphi(a) \frac{da}{a} 1_{(a \leq t)}\right], \text{ since } (I_t \leq a) = (t \geq T_a)$$

Now, the identity (10) (between the extreme terms)

(10)

After Rolle when we replace  $t$  by a generic stopping time. Hence, another use of the monotone class theorem yields:

$$(M) \quad E[F \varphi(I_\infty)] = E\left[\int_0^\infty \varphi(a) da \frac{F}{T_a} 1_{(T_a < \infty)}\right]$$

Finally, recalling that  $I_\infty = \bigcup_n I_n$  under  $P_n$  allows to deduce the identity (M) from (M1).

3. A concluding remark.

The statement of the theorem admits for a proof mixing initial enlargement with  $I_\infty$  and progressive enlargement with  $g$ .

In fact we have not proceeded exactly like this, the Proposition plays the role of the initial enlargement method and relies on a classical Gaussian relationship between  $F_2$  and  $P_2$ .

In conclusion, we like to present the above as an example (certainly among many of the same kind) of a mélange of enlargements techniques and Gaussian theorem.

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