

The distribution of Brownian and Bessel quantiles.

(A) /

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Second Edition

Introduction

Let $(B_t, t \geq 0)$ be a 1-dimensional Brownian motion starting from 0. The computation of the price of Asian options amounts to the computation of the following quantity:

$$(0.a) \quad C_T^{\nu}(k) \stackrel{\text{def}}{=} E \left[\left(\frac{1}{T} \int_0^T ds \exp(B_s + \nu s) - k \right)^+ \right],$$

for some given time T , and some reals ν and k . A number of results have been obtained about this topic (see H. Geman - M. Yor [1], [2], and [3], [4], in particular).

In November 1993, P. Embrechts suggested that, instead of considering, as in (0.a), the time-average of $\exp(B_s + \nu s)$ on the interval $[0, T]$, it would be interesting to study the median of this (random) function of $s (\leq T)$, and, more generally, the α -quantiles, for every $0 < \alpha < 1$. Precisely, to simplify the discussion, we may take, without loss of generality, $T=1$, and we define, for $0 \leq \alpha < 1$:

$$X_{\alpha}^{(\nu)} = \inf \left\{ x : \int_0^1 du \mathbb{1}(B_u + \nu u \leq x) > \alpha \right\}.$$

The main aim of this paper is to find the law of $X_{\alpha}^{(\nu)}$ explicitly, and to compute:

$$(0.b) \quad M^{\nu}(\alpha, k) \stackrel{\text{def}}{=} E \left[\left(X_{\alpha}^{(\nu)} - k \right)^+ \right].$$

Résumé de la version
Intermittent & cauchy-like performance,
+ ambrosius.

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Section 1 is devoted to this problem in the case $\gamma=0$, for which the formulae²⁾ are most explicit, whereas in section 2, the general case, with $\gamma \in \mathbb{R}$, is studied.

In Section 3, two variants of the preceding question are studied; the first variant concerns the study of the Brownian quantiles when the probability space $([0,1]; dt)$ is being replaced by $([0, \infty[; e^{-t} dt)$; thus, one is interested in:

$$Y_{\alpha}(\gamma) = \inf \left\{ x: \int_0^{\infty} dt e^{-t} 1_{(B_t + \gamma t \leq x)} > \alpha \right\};$$

the second variant concerns the Brownian quantiles over the probability space $([0, T_1], \frac{dt}{T_1})$, where $T_1 = \inf \{ t: B_t + \gamma t = 1 \}$; hence, one is interested in:

$$Z_{\alpha}(\gamma) = \inf \left\{ x: \frac{1}{T_1} \int_0^{T_1} dt 1_{(B_t + \gamma t \leq x)} > \alpha \right\}.$$

Finally, in Section 4, slightly different quantities are considered,

namely:

$$\eta_{\alpha}(\delta) = \inf \left\{ x: \int_0^{\infty} dt 1_{(R_t \leq x)} > \alpha \right\},$$

and

$$\mu_{\alpha}(\delta) = \inf \left\{ x: \int_0^{\infty} dt 1_{(|B_t| + \frac{2}{\delta} |t| \leq x)} > \alpha \right\}.$$

2. The general case: ν real.

Keeping with the notation in the introduction, we consider here:

$$X_\alpha^{(\nu)} = \inf \left\{ u : \int_0^1 du \mathbb{1}_{(B_u^{(\nu)} \leq x)} > \alpha \right\},$$

where $B_u^{(\nu)} = B_u + \nu u$, $u \geq 0$.

The symmetry relation (1.c) may be extended as follows:

$$(2.a) \quad X_\alpha^{(\nu)} \stackrel{(law)}{=} -X_{1-\alpha}^{(-\nu)}$$

Hence, to derive the distribution of $X_\alpha^{(\nu)}$, it suffices to compute $P(X_\alpha^{(\nu)} > x)$,

for every α , and $x \geq 0$.

With obvious notation, we remark, following the arguments developed in Section 1,

that:

$$\begin{aligned} (X_\alpha^{(\nu)} > x) &= (T_x^{(\nu)} + \int_{T_x^{(\nu)}}^1 dt \mathbb{1}_{(B_t^{(\nu)} < x)} < \alpha) \\ &\stackrel{(law)}{=} (T_x^{(\nu)} + \int_0^{1-T_x^{(\nu)}} dt \mathbb{1}_{(B_t^{(\nu)} < 0)} < \alpha). \end{aligned}$$

Hence, we obtain:

$$(2.b) \quad P(X_\alpha^{(\nu)} > x) = \int_0^\alpha dt \theta_x^{(\nu)}(t) P\left(\int_0^{1-t} du \mathbb{1}_{(B_u^{(\nu)} < 0)} < \alpha - t\right),$$

where:

$$\begin{aligned} \theta_x^{(\nu)}(t) dt &= P(T_x^{(\nu)} \in dt) = E\left[T_x \in dt; \exp\left(\nu x - \frac{\nu^2 t}{2}\right)\right] \\ &= \frac{dt x}{\sqrt{2\pi t^3}} \exp\left(-\frac{x^2}{2t} + \nu x - \frac{\nu^2 t}{2}\right) \end{aligned}$$

3. Two Variants.

The purpose of this section is to show how the arguments developed in Sections 1 and 2, although elementary, may yield results in somewhat different contexts.

(3.1) As a first variant, we shall consider:

$$Y_\alpha = \inf \left\{ x : \int_0^\infty dt e^{-t} 1_{(B_t < x)} > \alpha \right\} \quad (0 < \alpha < 1)$$

With the same notation as in (1.2) above, we find, for $x > 0$:

$$\begin{aligned} (Y_\alpha > x) &= \left(\int_0^{T_x} dt e^{-t} + \int_{T_x}^\infty e^{-t} 1_{(B_t < x)} dt < \alpha \right) \\ &= \left(\int_0^{T_x} dt e^{-t} + e^{-T_x} A < \alpha \right), \end{aligned}$$

where $A = \int_0^\infty du e^{-u} 1_{(B_u < 0)}$.

Hence, we obtain:

$$\begin{aligned} (Y_\alpha > x) &= (T_x < \log \frac{1-A}{1-\alpha}) \\ &\stackrel{\text{(law)}}{=} \left(\frac{x}{|N|} < \left(\log \frac{1-A}{1-\alpha} \right)^{1/2} \right). \end{aligned}$$

Similar arguments yield, for $x < 0$:

$$(Y_\alpha > x) \stackrel{\text{(law)}}{=} \left(\frac{-x}{|N|} > \left(\log \frac{A}{\alpha} \right)^{1/2} \right).$$

Consequently, we have obtained the following: The distribution of Y_α is characterized by:

Proposition:

Let $0 < \alpha < 1$.

$$P(Y_\alpha \in dx) = \begin{cases} \sqrt{\frac{2}{\pi}} E \left[\frac{1}{(\lambda_+)^{1/2}} \exp\left(-\frac{x^2}{2\lambda_+}\right) \right] & (x > 0) \\ \sqrt{\frac{2}{\pi}} E \left[\frac{1}{(\lambda_-)^{1/2}} \exp\left(-\frac{x^2}{2\lambda_-}\right) \right] & (x \leq 0), \end{cases}$$

where: $\lambda_+ = \log \frac{1-A}{1-\alpha}$, $\lambda_- = \log \frac{A}{\alpha}$, and $A = \int_0^\infty dt e^{-t} 1_{(B_t < 0)}$

(B) /

1)

The distribution of Brownian quantiles.

November 16th

Let (B_t) be a 1-dimensional Brownian motion starting from 0.

Define, for $0 < \alpha < 1$:

$$X_\alpha = \inf \left\{ x : \int_0^1 du \mathbb{1}(B_u < x) = \alpha \right\}.$$

I will show the following formula:

$$P(X_\alpha \in dx) = \begin{cases} \frac{\sqrt{2}}{\pi} E \left[g \leq \alpha ; \sqrt{\frac{1-g}{\alpha-g}} \exp \left(-\frac{x^2}{2} \left(\frac{1-g}{\alpha-g} \right) \right) \right] & (x > 0) \\ \frac{\sqrt{2}}{\pi} E \left[g \geq \alpha ; \frac{1}{\sqrt{1-\frac{\alpha}{g}}} \exp \left(-\frac{x^2}{2} \left(1 - \frac{\alpha}{g} \right) \right) \right] & (x \leq 0) \end{cases}$$

where g is arc sine distributed.

Proof:

I will use the following identity:

$$(X_\alpha > x) = \left(\int_0^1 ds \mathbb{1}(B_s < x) < \alpha \right).$$

(i) Assume $x > 0$; define $T_x = \inf \{ u : B_u = x \}$.

Then, we have: $P(X_\alpha > x) = P \left(T_x + \int_{T_x}^1 ds \mathbb{1}(B_s < x) < \alpha ; T_x < 1 \right)$

Since $(B_{T_x+u} - x; u \geq 0)$ is a Brownian motion independent of \mathcal{F}_{T_x} , we have:

$$P(X_\alpha > x) = P(T_x + (1-T_x)g < \alpha ; T_x < 1),$$

where g is arc sine distributed, and independent of T_x .

Therefore, we have:

$$\begin{aligned} P(X_\alpha > x) &= P(T_x(1-q) + q < \alpha; T_x < 1) \\ &= P(q < \alpha; x^2(1-q) < N^2; x^2 < N^2) \\ &= P(q < \alpha; x \leq \sqrt{\frac{\alpha-q}{1-q}} |N|). \end{aligned}$$

$$= \sqrt{\frac{2}{\pi}} E \left[q < \alpha; \sqrt{\frac{1-q}{\alpha-q}} \int_x^\infty dz \exp\left(-\frac{z^2}{2} \left(\frac{1-q}{\alpha-q}\right)\right) \right].$$

It now remains to differentiate with respect to x to obtain the formula for $x > 0$.

(ii) Assume $x < 0$.

Again, when I consider: $\left(\int_0^1 ds 1_{(B_s \leq x)} \leq \alpha\right)$,

I have to decompose with $(T_x \leq 1)$, and $(T_x > 1)$.

On $T_x > 1$, we have:

$$\int_0^1 ds 1_{(B_s \leq x)} = 0 \quad (< \alpha)$$

Therefore, we obtain:

$$\begin{aligned} P(X_\alpha > x) &= P(T_x > 1) + P(T_x \leq 1; \int_{T_x}^1 ds 1_{(B_s \leq x)} \leq \alpha) \\ &= P(T_x > 1) + P(T_x \leq 1; (1-T_x)q \leq \alpha) \\ &= P(|N| \leq |x|) + P(q \leq \alpha) P(|x| \leq |N|) + E \left[q \geq \alpha; 1 - \frac{\alpha}{q} \leq \frac{x^2}{N^2} \leq 1 \right] \end{aligned}$$

Therefore, writing $z = -x$, we get:

$$P(X_\alpha > x) = \sqrt{\frac{2}{\pi}} \int_0^z dy e^{-y^2/2} + P(q \leq \alpha) \sqrt{\frac{2}{\pi}} \int_z^\infty dy e^{-y^2/2} + E[(q \geq \alpha) \sqrt{\frac{2}{\pi}} \int_z^{\frac{z}{\sqrt{1-\frac{\alpha}{q}}}} dy \exp(-\frac{y^2}{2})]$$

Then, taking the derivative with respect to x on both sides, we get that the density, for $x < 0$, is:

$$\sqrt{\frac{2}{\pi}} E\left[q \geq \alpha; \frac{1}{\sqrt{1-\frac{\alpha}{q}}} \exp\left(-\frac{x^2}{2\left(1-\frac{\alpha}{q}\right)}\right)\right], \text{ which is the}$$

second formula \square .

Further simplifications: 1) Using the explicit distribution of q , we obtain the

following formula:

$$P(X_\alpha \in dx) = \begin{cases} \sqrt{\frac{2}{\pi}} E\left[\exp\left(-\frac{x^2}{2}\left(1 + \left(\frac{1-\alpha}{\alpha}\right)\frac{1}{q}\right)\right)\right] & (x \geq 0) \\ \sqrt{\frac{2}{\pi}} E\left[\exp\left(-\frac{x^2}{2}\left(1 + \left(\frac{\alpha}{1-\alpha}\right)\frac{1}{q}\right)\right)\right] & (x \leq 0) \end{cases}$$

It appears clearly with this formula that:

$-X_\alpha \stackrel{\text{(law)}}{=} X_{(1-\alpha)}$, as can be shown directly from the symmetry of BM.

2) The preceding formula may be simplified again; we introduce the following notation:

$$\frac{1-\alpha}{\alpha} = \beta^2 \quad (\beta > 0).$$

We will use the following

Lemma: $E \left[\exp \left(-\frac{y^2}{2} \left(\frac{1}{q} \right) \right) \right] = \Phi(|y|),$

where: $\Phi(z) = \sqrt{\frac{2}{\pi}} \int_z^{\infty} dx \exp\left(-\frac{x^2}{2}\right) = \mathbb{P}(|N| > z).$

Proof: We know that: $\frac{1}{q} \stackrel{\text{(law)}}{=} 1 + C^2$, where C is a standard

Cauchy variable. Then, we have:

$$\begin{aligned} E \left[\exp \left(-\frac{y^2}{2} \left(\frac{1}{q} \right) \right) \right] &= E \left[\exp \left(-\frac{y^2}{2} (1 + C^2) \right) \right] \\ &= \exp \left(-\frac{y^2}{2} \right) E \left[\exp (iy NC) \right] \\ &= \exp \left(-\frac{y^2}{2} \right) E \left[\exp (-|y| |N|) \right]. \\ &= \sqrt{\frac{2}{\pi}} \int_0^{\infty} dx \exp \left(-\frac{(x + |y|)^2}{2} \right) = \Phi(|y|). \end{aligned}$$

Now, we can write:

$$P(X_\alpha \in dx) = \begin{cases} \sqrt{\frac{2}{\pi}} \exp\left(-\frac{x^2}{2}\right) \Phi(\beta x) dx, & x \geq 0. \\ \sqrt{\frac{2}{\pi}} \exp\left(-\frac{x^2}{2}\right) \Phi\left(\frac{|x|}{\beta}\right) dx, & x \leq 0. \end{cases}$$

In the particular case where $\alpha = \frac{1}{2}$, we get: $\beta = 1$, and then the previous formula becomes:

$$\mathbb{P}\left(X_{\frac{1}{2}} \in dx\right) = \sqrt{\frac{2}{\pi}} \left(\exp\left(-\frac{x^2}{2}\right)\right) \mathbb{I}(|x|) dx, \quad (x \in \mathbb{R}).$$

Thus, we have: $\mathbb{P}\left(|X_{\frac{1}{2}}| \geq y\right) = \left(\mathbb{I}(y)\right)^2,$

which proves that: $|X_{\frac{1}{2}}| \stackrel{(kw)}{=} \inf(|N|, |N'|).$

(C) / Pour Thierry: Ceci est certainement la version la + proche de
Ce que fait ton élève —

1)

On the "true" Brownian quantiles

November, 16th, 1993

Let $(B_t, t \geq 0)$ be a 1-dimensional Brownian motion starting from 0.
For any fixed t , define, for $0 < \alpha < t$:

$$(1) \quad X_\alpha(t) = \inf \left\{ x : \int_0^t du \mathbf{1}_{(B_u \leq x)} = \alpha \right\}$$

We are interested in finding the distribution of $X_\alpha(1)$, for a given $\alpha < 1$.

We will take advantage of the scaling property of Brownian motion, from which we get:

$$(2) \quad X_\alpha(t) = \sqrt{t} X_{\frac{\alpha}{t}}(1)$$

Again, I keep t fixed. Since, $\alpha \rightarrow X_\alpha(t)$, $\alpha \leq t$, is the
inverse of the increasing process:

$$\alpha \rightarrow \int_0^t du \mathbf{1}_{(B_u \leq x)} \equiv \int_{-\infty}^x dy \ell_t^y,$$

we have, for any $f: \mathbb{R} \rightarrow \mathbb{R}_+$:

$$(3) \quad \int_{-\infty}^{\infty} dx \ell_t^x \exp\left(-\lambda \int_{-\infty}^x dy \ell_t^y\right) f(x) \\ = \int_0^t dx \exp(-\lambda x) f(X_\alpha(t)) \stackrel{(\text{law})}{=} t \int_0^1 d\beta \exp(-\lambda \beta t) f(\sqrt{t} X_\beta(1))$$

Moreover, if we assume f to have compact support, we get, by integration by parts, that (3) is equal to:

$$\frac{1}{\lambda} \int_{-\infty}^{\infty} dx \exp\left(-\lambda \int_{-\infty}^x dy \ell_t^y\right) f'(x).$$

Hence, this is equal, in law, to:

$$t \int_0^1 d\beta \exp(-\lambda\beta t) f(\sqrt{t} X_\beta(1)).$$

Now, the distribution of $(L_t^y; y \in \mathbb{R})$ is not so easy to manipulate - therefore, we randomize the time t , i.e.: we replace t by T_θ , an exponential random variable with parameter θ ; thus, we get:

$$(4) \quad E \left[T \int_0^1 d\beta \exp(-\lambda\beta T) f(\sqrt{T} X_\beta(1)) \right] \\ = \frac{1}{\lambda} \int_{-\infty}^{\infty} dx f'(x) E \left[\exp(-\lambda \int_0^T ds \mathbb{1}_{(B_s \leq x)}) \right]$$

$$[\quad \quad \quad !] .$$